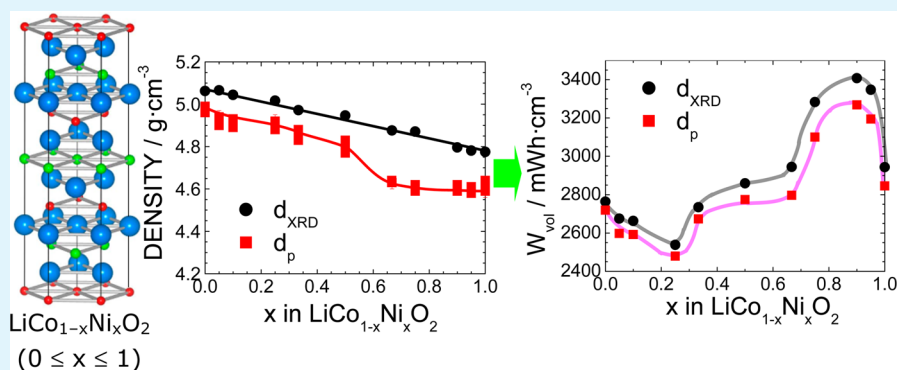


# Factors Affecting the Volumetric Energy Density of Lithium-Ion Battery Materials: Particle Density Measurements and Cross-Sectional Observations of Layered $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ with $0 \leq x \leq 1$

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**S** Supporting Information



**ABSTRACT:** Volumetric capacity  $Q_{\text{vol}}$  ( $\text{mAh cm}^{-3}$ ), more correctly, volumetric energy density  $W_{\text{vol}}$  ( $\text{mWh cm}^{-3}$ ), is a crucial property of lithium-ion battery (LIB) materials, because LIBs are devices that operate in a limited space. The actual value of  $W_{\text{vol}}$  ( $W_{\text{vol}}^{\text{act}}$ ) is currently limited to 40–60% of the maximum (theoretical) value of  $W_{\text{vol}}$  ( $W_{\text{vol}}^{\text{max}}$ ), for reasons that have not yet been fully clarified. Thus, to gain information that will enable an increase in  $W_{\text{vol}}^{\text{act}}$  such that it is closer to  $W_{\text{vol}}^{\text{max}}$ , systematic studies of the values for  $Q_{\text{vol}}$ ,  $W_{\text{vol}}$ , true density ( $d_{\text{XRD}}$ ), and particle density ( $d_{\text{p}}$ ) obtained using gas pycnometry were undertaken for  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $0 \leq x \leq 1$ . Here,  $d_{\text{p}}$  is the density that includes the volume of the closed pores in the particles, and consequently is less than  $d_{\text{XRD}}$ , which is determined by X-ray diffraction (XRD) measurement.  $d_{\text{XRD}}$  monotonically decreased from 5.062(1)  $\text{g cm}^{-3}$  for  $x = 0$  to 4.779(1)  $\text{g cm}^{-3}$  for  $x = 1$ , as expected. On the contrary,  $d_{\text{p}}$  decreased almost linearly from 4.98(2)  $\text{g cm}^{-3}$  for  $x = 0$  to 4.80(2)  $\text{g cm}^{-3}$  for  $x = 0.5$ , then suddenly dropped to 4.63(2)  $\text{g cm}^{-3}$  for  $x = 0.667$ , and finally leveled off to a constant value ( $\sim 4.6 \text{ g cm}^{-3}$ ) at larger values of  $x$ . The cross-sectional observations using a Focused Ion Beam system revealed that the significantly smaller values for  $d_{\text{p}}$  compared with those for  $d_{\text{XRD}}$ , particularly when  $x > 0.5$ , is due to the presence of closed pores in agglomerated secondary particles. This indicates that the closed pores in the secondary particles play an important role in determining the value of  $W_{\text{vol}}^{\text{act}}$  for LIBs. The formation of well-developed primary particles as a mean for increasing the value of  $d_{\text{p}}$  was also investigated.

**KEYWORDS:** lithium insertion material, high-energy density, XRD density, focused ion beam, particle morphology, closed pore

## 1. INTRODUCTION

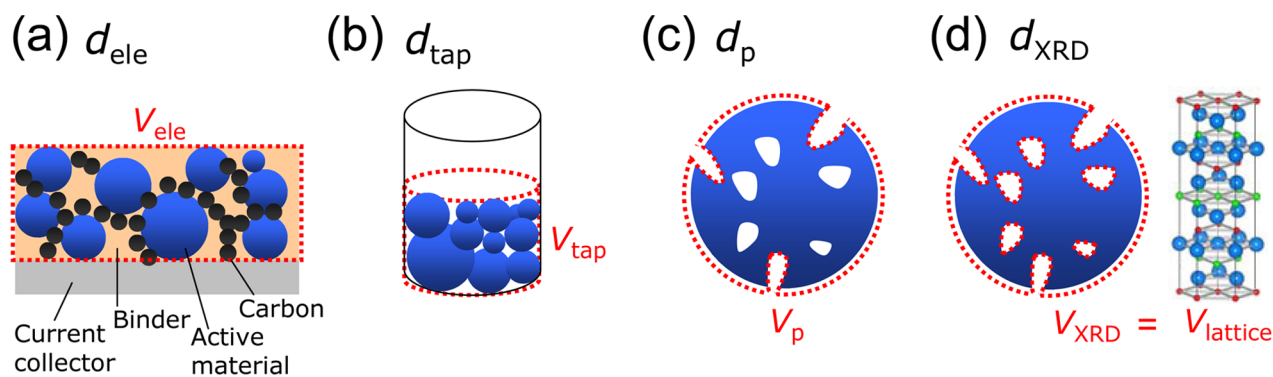
Lithium-ion batteries (LIBs) have received significant attention as storage devices for renewable energy because of their high energy density, high power performance, and long cycle life.<sup>1,2</sup> For high-energy density applications, lithium insertion materials with a layered structure, such as  $\text{LiCoO}_2$ ,<sup>3–7</sup>  $\text{LiNiO}_2$ ,<sup>8,9</sup> and  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  with  $0 < x < 1$ ,<sup>10–14</sup> have been employed as a positive electrode material. Here,  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  is a solid-solution compound between  $\text{LiCoO}_2$  and  $\text{LiNiO}_2$  over the entire  $x$  range, according to the monotonic change in the lattice parameters with  $x$ .<sup>10–14</sup> Recently, lithium insertion materials with a superlattice structure, such as  $\text{LiNi}_{1/2}\text{Mn}_{1/2}\text{O}_2$  and  $\text{LiNi}_{1/3}\text{Mn}_{1/3}\text{Co}_{1/3}\text{O}_2$ , have been intensively investigated because of their structural stability at high voltages above 4.2 V vs  $\text{Li}^+/\text{Li}$ .<sup>15</sup>

The energy density ( $W$ ) of a lithium insertion material is generally indicated by the gravimetric capacity ( $Q_{\text{gra}}$  ( $\text{mAh g}^{-1}$ )), most likely because the theoretical gravimetric capacity ( $Q_{\text{gra}}^{\text{theo}}$ ) can be easily calculated using the formula weight ( $M_w$ ) and Faraday constant. For instance,  $Q_{\text{gra}}^{\text{theo}} = 273.84 \text{ mAh g}^{-1}$  for  $\text{LiCoO}_2$ , assuming one-electron transfer per  $M_w$  of  $\text{LiCoO}_2$  ( $= 97.8730$ ) and 100% of Coulombic efficiency. However, because LIBs operate in a limited space, the volumetric capacity ( $Q_{\text{vol}}$  ( $\text{mAh cm}^{-3}$ )), rather than  $Q_{\text{gra}}$ , is important for evaluating lithium insertion materials. The volumetric energy density ( $W_{\text{vol}}$  ( $\text{mWh cm}^{-3}$ )) is given by

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**Figure 1.** Definition of several densities for lithium-ion battery materials: (a) density of an electrode ( $d_{\text{ele}}$ ), (b) tap density ( $d_{\text{tap}}$ ), (c) particle density ( $d_{\text{p}}$ ), and (d) density determined by X-ray diffraction (XRD) measurement ( $d_{\text{XRD}}$ ).  $d_{\text{ele}}$  is the density of the active material (lithium insertion material) in the electrode volume ( $V_{\text{ele}}$ ).  $V_{\text{ele}}$  contains the volumes for the conductive carbon and binder.  $d_{\text{tap}}$  is calculated using the weight of the active material and tapped volume ( $V_{\text{tap}}$ ) in a container.  $d_{\text{ele}}$  and  $d_{\text{tap}}$  are strongly dependent on the experimental conditions, such as the particle size, particle shape, particle distribution of the active material,  $V_{\text{tap}}$ , and the compaction process. On the other hand,  $d_{\text{p}}$  and  $d_{\text{XRD}}$  are well-defined values.  $d_{\text{XRD}}$  is determined using the formula weight and lattice volume ( $V_{\text{lattice}} = V_{\text{XRD}}$ ), whereas  $d_{\text{p}}$  is calculated using the volume ( $V_{\text{p}}$ ) that includes the closed pores in a particle, which do not have contact with the particle surface. Thus, the magnitudes of the densities are in the order  $d_{\text{ele}} < d_{\text{tap}} < d_{\text{p}} < d_{\text{XRD}}$ .

$$W_{\text{vol}} = \int_0^{Q_{\text{vol}}} V dQ = V_{\text{ave}} Q_{\text{vol}} \quad (1)$$

where  $V_{\text{ave}}$  is the average operating voltage.

The maximum volumetric capacity ( $Q_{\text{vol}}^{\text{max}}$ ) of a lithium insertion material is obtained using  $Q_{\text{gra}}$  and the true density, i.e., the density determined by powder X-ray diffraction (XRD) measurement ( $d_{\text{XRD}}$  ( $\text{g cm}^{-3}$ )).<sup>16</sup> Furthermore, since  $V_{\text{ave}}$  is derived by ab initio calculations,<sup>17</sup> the maximum volumetric energy density ( $W_{\text{vol}}^{\text{max}}$ ) can be determined. Nevertheless, it remains difficult to predict the actual volumetric capacity ( $Q_{\text{vol}}^{\text{act}}$ ) or actual energy density ( $W_{\text{vol}}^{\text{act}}$ ), because the actual density in the electrode ( $d_{\text{ele}}$ ) is strongly affected by the particle shape, particle size, and particle distribution of the lithium insertion material. Moreover,  $d_{\text{ele}}$  is influenced by the amounts and/or types of conducting carbon and binder, which are invariably used in LIB electrodes. Although  $Q_{\text{vol}}^{\text{act}}$  ( $W_{\text{vol}}^{\text{act}}$ ) is empirically known to be 40–60% of  $Q_{\text{vol}}^{\text{max}}$  ( $W_{\text{vol}}^{\text{max}}$ ),<sup>2</sup> factors affecting the difference between  $d_{\text{XRD}}$  and  $d_{\text{ele}}$  should be clarified in order to bring  $Q_{\text{vol}}^{\text{act}}$  ( $W_{\text{vol}}^{\text{act}}$ ) close to  $Q_{\text{vol}}^{\text{max}}$  ( $W_{\text{vol}}^{\text{max}}$ ).

Figure 1 schematically depicts several possible density definitions for LIB materials: (a)  $d_{\text{ele}}$ , (b) tap density ( $d_{\text{tap}}$ ), (c) particle density ( $d_{\text{p}}$ ), and (d)  $d_{\text{XRD}}$ . As described above,  $d_{\text{ele}}$  is calculated using the volume of the electrode ( $V_{\text{ele}}$ ) and the weight of the active material (lithium insertion material) in the electrode.  $d_{\text{tap}}$  is the density when the active material is tapped in a container with a volume  $V_{\text{tap}}$ .  $d_{\text{p}}$  is very close to  $d_{\text{XRD}}$ , but it is calculated using the volume ( $V_{\text{p}}$ ) that includes the “closed” pores in a particle. Here, closed refers to pores that do not have contact with the particle surface (see Figure 1c).  $d_{\text{XRD}}$  is calculated by  $M_{\text{w}}$  and lattice volume  $V_{\text{lattice}} (= V_{\text{XRD}})$ . The magnitudes of the different densities are in the order  $d_{\text{ele}} < d_{\text{tap}} < d_{\text{p}} < d_{\text{XRD}}$ , where the latter two densities are relatively well-defined values. Thus, as a first step, it is crucial to understand the factors affecting the difference in the values for  $d_{\text{p}}$  and  $d_{\text{XRD}}$ , in order to increase  $W_{\text{vol}}^{\text{act}}$ . For this purpose, a gas (expansion) pycnometer, which is used for measuring the density of polycrystalline powders, single-crystals, and liquids,<sup>18</sup> was employed. A liquid pycnometer was previously used to investigate the oxygen deficiency of  $\text{LiM}_x\text{Mn}_{2-x}\text{O}_{4-\delta}$  spinels with  $M = \text{Mg}, \text{Ni}, \text{etc.}$ <sup>19,20</sup> and to study the  $d_{\text{p}}$  of  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$ .<sup>21</sup> However, to the best of our knowl-

edge, systematic  $d_{\text{p}}$  measurements associated with  $Q_{\text{vol}}$  ( $W_{\text{vol}}$ ) have never been performed for lithium insertion materials.

Herein, the results of an investigation of a series of  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  materials with  $0 \leq x \leq 1$ , which are compounds commonly used as positive electrode materials for high-energy density LIBs,<sup>14</sup> are described. The  $x$  dependences of  $d_{\text{XRD}}$ ,  $d_{\text{p}}$ ,  $Q_{\text{vol}}$ , and  $W_{\text{vol}}$  were examined in order to obtain essential information required for designing advanced high-energy density LIBs. From the viewpoints of materials science, such density analyses would be important for in-depth understanding of electronic and thermal transport properties, fluid mechanics, and adsorption property of lithium transition metal oxides. This is because  $d_{\text{p}}$  is one of physical properties for agglomerated (macroscopic) powders, which is usually contrasted with a single (microscopic) particle based on statistical mechanics.<sup>18</sup> As a result, it was found that the value of  $d_{\text{p}}$  is significantly less than that of  $d_{\text{XRD}}$ , particularly for  $x > 0.5$ , because of the presence of closed pores in agglomerated secondary particles. In addition, methods for increasing the value of  $d_{\text{p}}$  for the composition with  $x = 0.75$  were investigated.

## 2. EXPERIMENTAL SECTION

**2.1. Sample Preparation.** Polycrystalline  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $x = 0, 0.05, 0.1, 0.25, 0.333, 0.5, 0.667, 0.75, 0.9, 0.95,$  and 1 were prepared using a solid-state reaction technique as reported previously.<sup>13,14</sup> First,  $\text{Co}_3\text{O}_4$  and  $\text{NiO}$  powders were prepared from  $\text{CoO}$  (Kojundo Chemical Lab. Co. Ltd.) and basic nickel carbonate (Wako Pure Chemical Industries, Ltd.), respectively, by heating at 750 °C for 12 h in air. For the samples with  $x < 0.75$ ,  $\text{Li}_2\text{CO}_3$  (Wako Pure Chemical Industries, Ltd.),  $\text{Co}_3\text{O}_4$ , and  $\text{NiO}$  were used as the starting materials. Each reaction mixture was well mixed using a mortar and pestle, and then pressed into a pellet with a diameter of 23 mm and a thickness of  $\sim 5$  mm. The pellet was subsequently heated in air for 12 h at 900 °C for  $x \leq 0.1$ , 850 °C for  $x = 0.25$  and 0.333, and 800 °C for  $x = 0.5$  and 0.667, respectively. For the samples with  $x \geq 0.75$ ,  $\text{LiOH} \cdot \text{H}_2\text{O}$  (Wako Pure Chemical Industries, Ltd.),  $\text{Co}_3\text{O}_4$ , and  $\text{NiO}$  were used as the starting materials. Each pellet formed from the different reaction mixtures was preheated in an oxygen flow at 650 °C for 12 h, then ground and repressed into a pellet, and finally subjected to an oxygen flow at 750 °C for 12 h.

For the composition with  $x = 0.75$ , the samples were also prepared by a flux method using  $\text{NaCl}$  or  $\text{KCl}$  (Wako Pure Chemical Industries, Ltd.) as the reaction medium. The  $\text{LiOH} \cdot \text{H}_2\text{O}/\text{Co}_3\text{O}_4/\text{NiO}$  mixture

and the NaCl or KCl medium were ground together using a mortar and pestle, and then heated at 750 °C in air for 18 h. The ratio of the reaction mixture to the medium was 1:3. The obtained powders were washed with distilled water and dried at 200 °C for 24 h in a vacuum.

## 2.2. Structural, Morphology, and Electrochemical Analyses.

The crystal structure of the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples was examined by powder XRD measurements with copper  $K\alpha$  radiation (RINT-2200, Rigaku Co. Ltd.) and iron  $K\alpha$  radiation (D8 ADVANCE, Bruker AXS, Inc.). The lattice parameters in the hexagonal setting, i.e.,  $a_h$  and  $c_h$  were determined by a Rietveld analysis with RIETAN2000 software.<sup>22</sup> Particle morphologies were investigated by scanning electron microscopy (SEM) analyses (S-3600N, Hitachi High-Technologies Co. Ltd.). Electrochemical properties were investigated in a non-aqueous lithium cell, as reported previously.<sup>6,7,13,14</sup> Polyvinylidene fluoride (PVdF) dissolved in *N*-methyl-2-pyrrolidone (NMP) solution was used as the binder when preparing the mixed electrodes. A black viscous slurry consisting of 88 wt %  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ , 6 wt % acetylene black (AB), and 6 wt % PVdF was cast on an aluminum foil with a blade to form a mixed electrode ( $\varphi = 16$  mm), which was then dried under a vacuum at 120 °C for 12 h in order to evaporate NMP. A lithium metal sheet pressed on a stainless steel plate ( $\varphi = 19$  mm) was used as the counter electrode. Two sheets of porous polyethylene membrane (TonenGeneral Sekiyu K. K.) were used as the separator. The electrolyte was composed of 1 M  $\text{LiPF}_6$  dissolved in an ethylene carbonate/dimethyl carbonate (1/1 volume ratio) solution (Kishida Chemical Co., Ltd.). The cell was assembled in an argon-filled glovebox. The charge and discharge tests were performed in the voltage range between 3.0 and 4.2 V at a current density of 0.15 mA  $\text{cm}^{-2}$  and 25 °C.

**2.3. Measurements of  $d_p$ .** The  $d_p$  values were determined by gas pycnometry (Ultrapycnometer 1000, Quantachrome Instruments Inc.), which is based on the Archimedes' displacement principle.<sup>18</sup> With gas pycnometry, volume displacement is not measured directly but is determined using the equation of state, i.e.,  $PV = nRT$ . First, the volume of the sample cell ( $V_c$ ) and the volume of the added space ( $V_{\text{add}}$ ) were calibrated using calibration spheres certified by the National Institute of Standards and Technology (NIST, 01500-NISTSM). A schematic illustration of the pycnometer system is shown in Figure S1 in the Supporting Information. Helium gas was used as the medium because it behaves as an ideal gas and penetrates pores smaller than 2 Å. Thus, the equation of state for the sample cell is represented by

$$P_a V_c = n_c RT_a \quad (2)$$

and for the added space, it is as follows

$$P_a V_{\text{add}} = n_{\text{add}} RT_a \quad (3)$$

where  $P_a$  is the ambient pressure ( $\sim 101$  kPa),  $n_c$  and  $n_{\text{add}}$  are the moles of helium gas in the sample cell and the added space, respectively,  $R$  is the gas constant, and  $T$  is the ambient temperature ( $= 25 \pm 0.5$  °C). When a  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  sample is placed in the sample cell, the equation of state for the sample cell can be described as

$$P_a (V_c - V_s) = n_1 RT_a \quad (4)$$

where  $V_s$  is the volume of the sample, and  $n_1$  is the moles of remaining the helium gas. The weight of the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  ( $w_s$ ) samples in this study was  $\sim 1.8$  g for all measurements. The sample cell was then pressurized to  $P_2$  ( $\sim 138$  kPa)

$$P_2 (V_c - V_s) = n_2 RT_a \quad (5)$$

where  $n_2$  is the remaining moles of helium gas at  $P_2$ . When the valve between  $V_c$  and  $V_{\text{add}}$  was opened, the pressure dropped to an equilibrium value of  $P_3$

$$P_3 (V_c - V_s + V_{\text{add}}) = n_2 RT_a + n_{\text{add}} RT_a \quad (6)$$

Using eqs 3 and 5,  $V_s$  was determined as

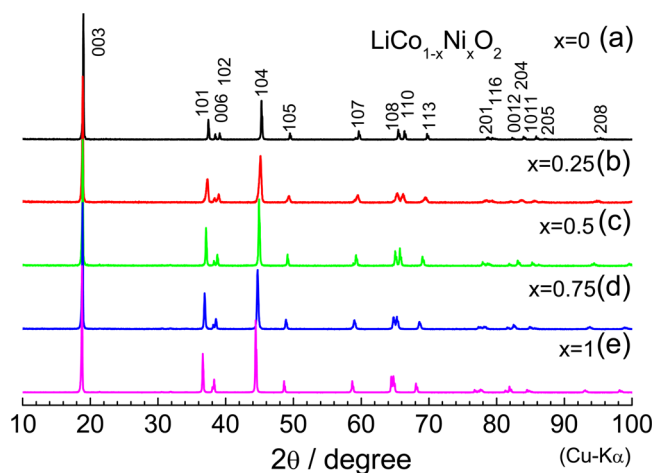
$$V_s = V_c + \frac{V_a}{\left(1 - \frac{P_2 - P_a}{P_3 - P_a}\right)} = V_c + \frac{V_a}{\left(1 - \frac{P_2}{P_3}\right)} \quad (7)$$

and finally, the value for  $d_p$  ( $= w_s/V_s$ ) was obtained. Therefore, it is possible to determine the value of  $d_p$  with just one measurement. However, in order to obtain a reliable  $d_p$  value, the measurement was repeated 100 times for each sample. That is, after one  $d_p$  measurement, the sample cell was repressurized to  $P_2$ , the valve between  $V_c$  and  $V_{\text{add}}$  was opened again, and  $V_s$  ( $d_p$ ) was calculated. The averaged  $d_s$  value for each sample was determined using the data for the last 50 measurements. The  $d_p$  value for  $x = 0$  as a function of each measurement is shown in Figure S2 of the Supporting Information. In addition, at least three independent runs were performed for each  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  sample in order to confirm the reproducibility of the obtained  $d_p$  values.

**2.4. Cross-Sectional Observations.** Cross-sectional views of the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  electrodes were observed using a focused ion beam (FIB) system with FB-2000 (Hitachi High-Tech Science Systems Co.). The electrodes consisted of 88 wt %  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ , 6 wt % AB, and 6 wt % PVdF. Three ranges of  $\text{Ga}^+$  ion currents were applied in order to cut the electrodes: 20 nA for 100 min, 2.6 nA for 30 min, and 0.62 nA for 10 min. The dimensions of the cutting zone were approximately 30  $\mu\text{m}$  in width, 50  $\mu\text{m}$  in height, and 5  $\mu\text{m}$  in depth. After cutting, each electrode was tilted  $\sim 90^\circ$ , and then the cross section of the electrode was observed using  $\text{Ga}^+$ -induced SEM. During SEM observation, the  $\text{Ga}^+$ -ion current was maintained at 0.15 nA in order to avoid possible beam damage. The  $\text{Ga}^+$  ions were accelerated at 30 kV under a vacuum of  $\sim 1 \times 10^{-4}$  Pa.

## 3. RESULTS AND DISCUSSION

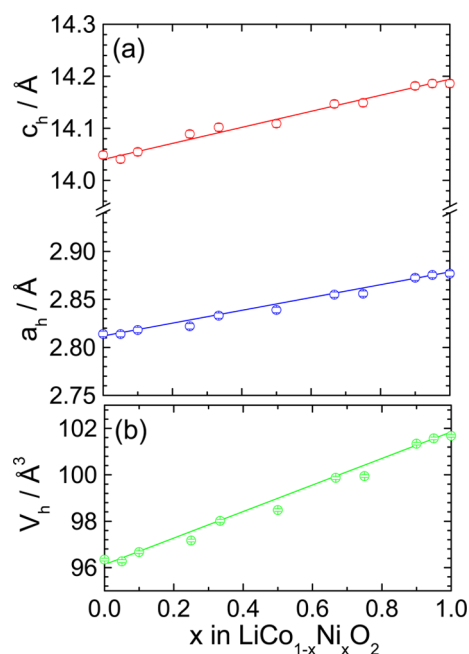
**3.1. Crystal Structure.** Figure 2 shows the XRD patterns for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with (a)  $x = 0$ , (b)  $x = 0.25$ , (c)



**Figure 2.** XRD patterns of the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with (a)  $x = 0$ , (b)  $x = 0.25$ , (c)  $x = 0.5$ , (d)  $x = 0.75$ , and (e)  $x = 1$ .

(d)  $x = 0.75$ , and (e)  $x = 1$ . The XRD patterns for the other  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples, i.e.,  $x = 0.05, 0.1, 0.333, 0.667, 0.9$ , and  $0.95$ , are shown in Figure S3 of the Supporting Information. The XRD patterns with the  $x \geq 0.75$  samples were identified as a single-phase of a layered structure with the  $R\bar{3}m$  space group, in which the  $\text{Li}^+$  ions occupy the octahedral  $3b$  site, whereas the  $\text{Co}^{3+}$  ( $\text{Ni}^{3+}$ ) ions occupy the octahedral  $3a$  site. For the samples with  $x \geq 0.9$ , the Rietveld analysis with RIETAN2000<sup>22</sup> clarified the presence of a small amount of Ni ions in the  $3b$  (Li) site; more specifically,  $z$  in  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  is 0.01(1) for  $x = 0.9$ , 0.02(1) for  $x = 0.95$ , and 0.03(1) for  $x = 1$ . Results of the Rietveld

analysis are shown in Figure S4 of the Supporting Information. Figure 3a shows the lattice parameters for the  $a_h$ - and  $c_h$ -axes,



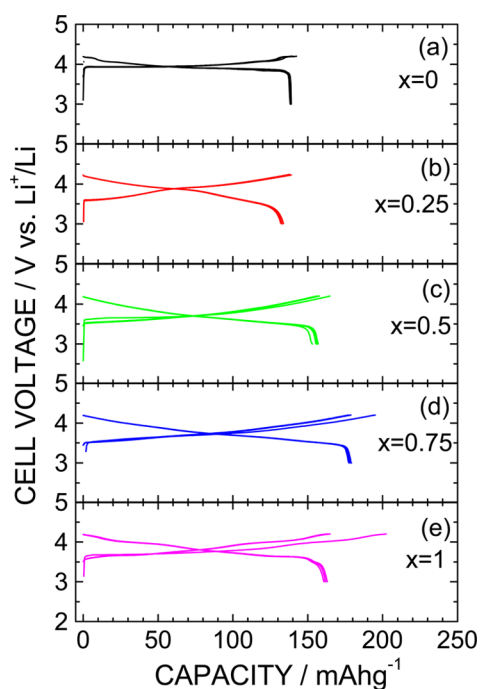
**Figure 3.** (a) Lattice parameters for the  $a_h$ - and  $c_h$ -axes and (b) unit-cell volume  $V_h$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $0 \leq x \leq 1$ . The values of  $a_h$  and  $c_h$  values were determined by the Rietveld analyses.

respectively, as a function of  $x$  in  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ . Both  $a_h$  and  $c_h$  increases with the value of  $x$ , indicating that the Co and Ni ions are homogeneously distributed in the Co/Ni layers over the entire range of  $x$  values. The increase in  $a_h$  and  $c_h$  with  $x$  is understood by the increase in the ionic radius ( $r$ ) resulting when  $\text{Co}^{3+}$  ions are replaced with  $\text{Ni}^{3+}$  ions. That is, when the coordination number is six for both ions,  $r = 0.53$  Å for the  $\text{Co}^{3+}$  ions in a low-spin state ( $t_{2g}^6 e_g^0$ ), whereas  $r = 0.56$  Å for the  $\text{Ni}^{3+}$  ions in a low-spin state ( $t_{2g}^6 e_g^1$ ).<sup>2,3</sup>

Due to the linear relationship between  $a_h$  ( $c_h$ ) and  $x$ , the variation of the unit cell volume  $V_h$  also exhibits a monotonic change with  $x$  (Figure 3b). Therefore, the structural parameters of the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples prepared in this study, which are

summarized in Table 1, are consistent with the previous results reported for  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ .<sup>10–14</sup>

**3.2. Electrochemical Properties.** Figure 4 presents the charge and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with



**Figure 4.** Charge and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with (a)  $x = 0$ , (b)  $x = 0.25$ , (c)  $x = 0.5$ , (d)  $x = 0.75$ , and (e)  $x = 1$ . The cells were operated in the voltage range between 3.0 and 4.2 V at a current density of  $0.15 \text{ mA cm}^{-2}$  and  $25^\circ\text{C}$ .

(a)  $x = 0$ , (b)  $x = 0.25$ , (c)  $x = 0.5$ , (d)  $x = 0.75$ , and (e)  $x = 1$ . The charge and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with  $x = 0.05, 0.1, 0.333, 0.667, 0.9$ , and  $0.95$  can be seen in Figure S5 of the Supporting Information. All of the cells were operated in the voltage range between 3.0 and 4.2 V at a current density of  $0.15 \text{ mA cm}^{-2}$  and  $25^\circ\text{C}$ . The charge cutoff voltage was restricted to 4.2 V, in order to avoid capacity fading due to the inherent structural phase transitions of  $\text{LiCoO}_2$ <sup>4–7</sup> and  $\text{LiNiO}_2$ ,<sup>8,9</sup> and/or degradation of the electrolytes at high voltages. For the  $x = 0$  sample, the cell voltage ( $V_{\text{cell}}$ ) initially

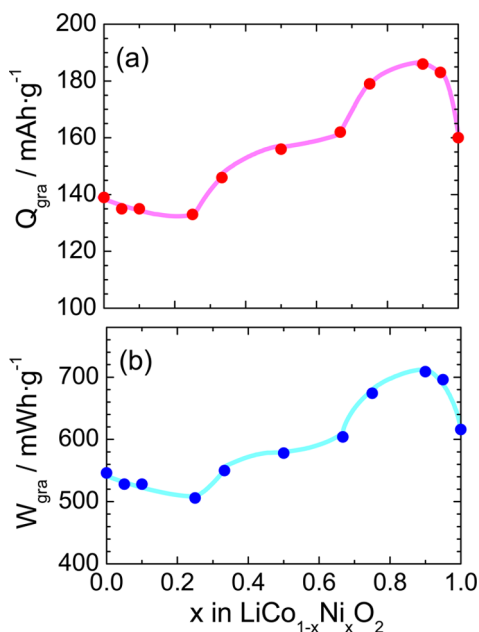
**Table 1.** Structural Parameters of  $a_h$ ,  $c_h$ , and  $V_h$ , Densities of  $d_{\text{XRD}}$  and  $d_p$ , Capacities of  $Q_{\text{gra}}$  and  $Q_{\text{vol}}$ , and Energy Densities of  $W_{\text{gra}}$  and  $W_{\text{vol}}$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  Samples with  $0 \leq x \leq 1$

$x$	structural parameters		density			capacity and energy density		
	lattice parameters (Å)	lattice volume (Å <sup>3</sup> )	$d_{\text{XRD}}$ (g cm <sup>-3</sup> )	$d_p$ (g cm <sup>-3</sup> )	$Q_{\text{gra}}$ (mAh g <sup>-1</sup> )	$W_{\text{gra}}$ (mWh g <sup>-1</sup> )	$Q_{\text{vol}}^a$ (mAh cm <sup>-3</sup> )	$W_{\text{vol}}^a$ (mWh cm <sup>-3</sup> )
0	$a_h = 2.814(1), c_h = 14.05(1)$	96.34(1)	5.062(1)	4.98(2)	139	546	692	2720
0.05	$a_h = 2.814(1), c_h = 14.04(1)$	96.27(1)	5.066(1)	4.92(2)	135	528	664	2598
0.1	$a_h = 2.818(1), c_h = 14.05(1)$	96.67(1)	5.044(1)	4.91(2)	135	528	663	2592
0.25	$a_h = 2.822(1), c_h = 14.09(1)$	97.17(1)	5.017(1)	4.90(2)	133	506	652	2479
0.333	$a_h = 2.833(1), c_h = 14.10(1)$	98.02(1)	4.972(1)	4.86(2)	146	550	710	2673
0.5	$a_h = 2.839(1), c_h = 14.11(1)$	98.48(1)	4.947(1)	4.80(2)	156	578	749	2774
0.667	$a_h = 2.855(1), c_h = 14.15(1)$	99.86(1)	4.876(1)	4.63(2)	162	604	751	2797
0.75	$a_h = 2.856(1), c_h = 14.15(1)$	99.95(1)	4.871(1)	4.60(2)	179	674	823	3100
0.9	$a_h = 2.873(1), c_h = 14.18(1)$	101.3(1)	4.806(1)	4.61(2)	186	709	858	3268
0.95	$a_h = 2.875(1), c_h = 14.19(1)$	101.6(1)	4.809(1)	4.59(2)	183	696	840	3194
1	$a_h = 2.877(1), c_h = 14.19(1)$	101.7(1)	4.779(1)	4.62(2)	161	616	739	2846

<sup>a</sup>The values of  $Q_{\text{vol}}$  and  $W_{\text{vol}}$  were calculated using the  $d_p$  value.

increases rapidly from  $\sim 3.0$  V, and then keeps nearly constant at  $\sim 3.9$  V up to  $Q_{\text{gra}} \approx 70 \text{ mAh g}^{-1}$ . The constant operating voltage of  $\sim 3.9$  V corresponds to the formation of two  $R\bar{3}m$  phases, which is characteristics for stoichiometric  $\text{LiCoO}_2$ .<sup>4–7</sup> As  $x$  increases from 0 to 0.25, the operating voltage at the beginning of the charge reaction decreases to  $\sim 3.6$  V (see Figures 4b and Figure S5 in the Supporting Information), indicating that the  $\text{Ni}^{3+}$  ions are oxidized before the  $\text{Co}^{3+}$  ions.<sup>12</sup> For the  $0.333 \leq x \leq 0.75$  samples, the  $V_{\text{cell}}$  almost monotonically increases with  $Q_{\text{gra}}$ , suggesting a one-phase reaction; in other words, no structural phase change occurs during the charge and discharge reactions. For the  $x \geq 0.9$  samples, however, plateaus are observed at around 3.6, 4.0, and 4.2 V. The former two plateaus suggest the formation and annihilation of a monoclinic ( $C2/m$ ) phase, respectively, whereas the last plateau indicates the formation of two  $R\bar{3}m$  phases in delithiated  $\text{Li}_y\text{NiO}_2$  with  $y < 0.5$ .<sup>8,9,13,14</sup>

In Figure 5a, the values for  $Q_{\text{gra}}$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples are plotted as a function of  $x$ . The values for  $Q_{\text{gra}}$  in



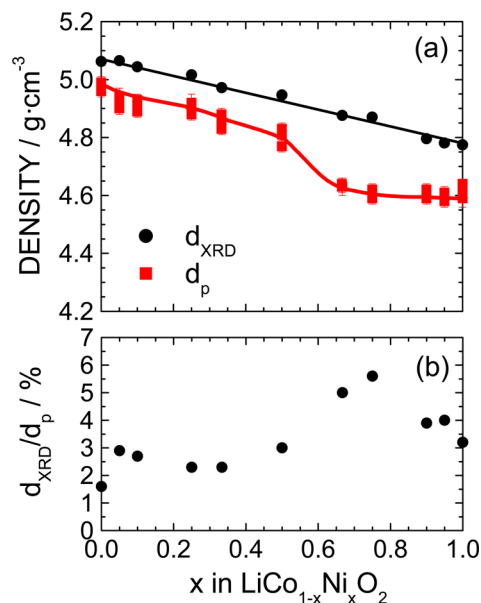
**Figure 5.** (a) Gravimetric discharge capacity ( $Q_{\text{gra}}$ ) and (b) gravimetric energy density ( $W_{\text{gra}}$ ) as a function of  $x$  in  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ .  $Q_{\text{gra}}$  is the discharge capacity for the 2nd cycle shown in Figure 4 and Figure S5 in the Supporting Information.  $W_{\text{gra}}$  was obtained by multiplying  $Q_{\text{gra}}$  by the average operating voltage ( $V_{\text{ave}}$ ).

the figure are the discharge capacities for the second cycle, as shown in Figure 4 and Figure S5 in the Supporting Information. The  $x$  dependence of  $Q_{\text{gra}}$  can be divided into three distinct regions; (i)  $0 \leq x \leq 0.25$ , (ii)  $0.25 < x < 0.75$ , and (iii)  $x \geq 0.75$ . In region (i),  $Q_{\text{gra}}$  slightly decreases from  $139 \text{ mAh g}^{-1}$  to  $133 \text{ mAh g}^{-1}$ , while in region (ii),  $Q_{\text{gra}}$  increases with a change in the slope ( $dQ_{\text{gra}}/dx$ ) and approaches  $162 \text{ mAh g}^{-1}$  for  $x = 0.667$ . Finally, in region (iii),  $Q_{\text{gra}}$  rapidly increases from  $x = 0.75$  and reaches a maximum ( $= 186 \text{ mAh g}^{-1}$ ) for the  $x = 0.9$  sample, and then declines to  $161 \text{ mAh g}^{-1}$  when  $x = 1$ . In our previous magnetic and electrochemical studies of  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  with  $x \geq 0.75$ , a similar  $x$  dependence was observed in region (iii), although in that case,  $Q_{\text{gra}} = 206 \text{ mAh g}^{-1}$  for  $x = 0.9$ ,  $Q_{\text{gra}} = 217 \text{ mAh g}^{-1}$  for  $x = 0.95$ , and  $Q_{\text{gra}} = 180 \text{ mAh g}^{-1}$  for  $x = 1$ .<sup>14</sup> The difference in the values for  $Q_{\text{gra}}$

obtained in this and previous studies is attributed to the difference in the  $z$  amounts in  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$ ,<sup>14</sup> as described below. The increase in  $Q_{\text{gra}}$  with  $x$  can be explained by a decrease in the redox potentials below 4.2 V with  $x$ .<sup>24</sup> Although there are three major redox potentials at 3.91, 4.18, and 4.50 V for  $x = 0$ , the former two redox potentials are decreased to 3.58 and 4.05 V when  $x = 0.5$  because of the formation of the  $\text{Ni}^{3+}/\text{Ni}^{4+}$  redox couple.<sup>24</sup> In region (iii),  $z$  also contributes to  $Q_{\text{gra}}$ . Because the presence of  $z$  disturbs the diffusion of  $\text{Li}^+$  ions in the Li layer, electrochemical properties, such as  $Q_{\text{gra}}$  and irreversible capacity, strongly depend on the quantity of  $z$ , as reported previously.<sup>9–14</sup> Indeed, as clarified by the Rietveld analysis (Section 3.1), the values of  $z$  for the  $x = 0.9$  and  $0.95$  samples are smaller than that for the  $x = 1$  sample. Thus, the  $\text{Ni}^{3+}/\text{Ni}^{4+}$  redox couple and the quantity of  $z$  in  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  mainly determine the values of  $Q_{\text{gra}}$ , which was measured in the voltage range between 3.0 and 4.2 V.

The  $W_{\text{gra}}$  value is then obtained by multiplying  $Q_{\text{gra}}$  and  $V_{\text{ave}}$ . As shown in Figure S6 in the Supporting Information, the value of  $V_{\text{ave}}$  is minimal for the  $x = 0.5$  sample. That is,  $V_{\text{ave}} = 3.93$ , 3.71, and 3.80 V for the  $x = 0, 0.5$ , and 1 samples, respectively. However, since the difference in  $V_{\text{ave}}$  for  $x = 0$  and 0.5 is within  $\sim 6\%$ , it can be concluded that  $V_{\text{ave}}$  does not strongly affect  $W_{\text{gra}}$  for  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ . In fact, the  $x$  dependence of  $W_{\text{gra}}$  is quite similar to that of  $Q_{\text{gra}}$  (Figure 5b); as  $x$  increases from 0 to 0.25,  $W_{\text{gra}}$  slightly decreases from  $546 \text{ mWh g}^{-1}$  to  $506 \text{ mWh g}^{-1}$ , then increases with further increasing  $x$ , and finally reaches a maximum ( $= 709 \text{ mWh g}^{-1}$ ) for  $x = 0.9$ . The values of  $Q_{\text{gra}}$  and  $W_{\text{gra}}$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples are listed in Table 1.

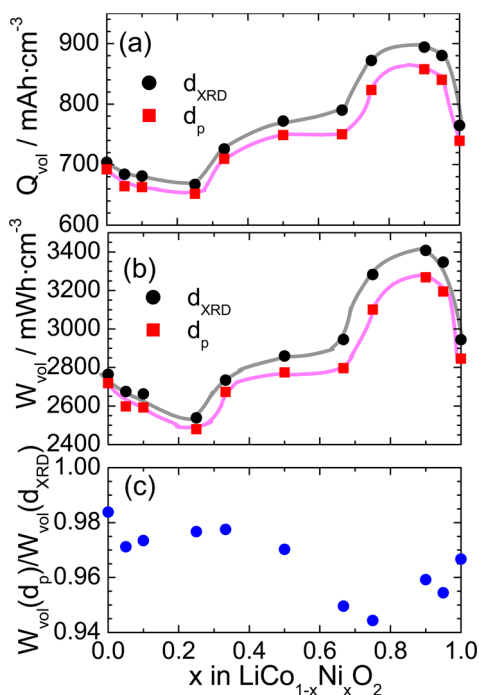
**3.3.  $d_p$  Measurements.** Figure 6a shows the values of  $d_p$  and  $d_{\text{XRD}}$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples. The  $d_p$  and  $d_{\text{XRD}}$  values are also summarized in Table 1, where  $d_p$  is the average of three independent runs. The  $d_{\text{XRD}}$  value almost linearly decreases with increasing  $x$ ;  $d_{\text{XRD}} = 5.062(1)$ ,  $4.947(1)$ , and  $4.779(1) \text{ g cm}^{-3}$  for the  $x = 0, 0.5$ , and 1 samples, respectively. The



**Figure 6.** (a) Particle density ( $d_p$ ) and true density ( $d_{\text{XRD}}$ ) as a function of  $x$  in  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ .  $d_{\text{XRD}}$  was determined by the XRD measurements. (b) The difference between  $d_p$  and  $d_{\text{XRD}}$ , i.e., the  $d_{\text{XRD}}/d_p$  ratio.

decrease in  $d_{\text{XRD}}$  with  $x$  is due to the increase in  $V_h$  with  $x$  (see Figure 3b), because  $M_w$  for  $\text{LiCoO}_2$  ( $= 97.8730$ ) is close to that for  $\text{LiNiO}_2$  ( $= 97.6332$ ). On the contrary, the  $x$  dependence of  $d_p$  is not monotonic. As  $x$  increases from 0 to 0.5,  $d_p$  almost decreases linearly from  $4.98(2) \text{ g cm}^{-3}$  to  $4.80(2) \text{ g cm}^{-3}$ , then suddenly drops to  $4.63(2) \text{ g cm}^{-3}$  for  $x = 0.667$ , and finally levels off at a constant value ( $\sim 4.6 \text{ g cm}^{-3}$ ) as  $x$  increases further. As can be seen in Figure 6b, the difference between  $d_{\text{XRD}}$  and  $d_p$ , i.e., the  $d_{\text{XRD}}/d_p$  ratio, has a minimum value at  $x = 0$  ( $= 1.6\%$ ). Then, the  $d_{\text{XRD}}/d_p$  is almost constant ( $\sim 3.0\%$ ) until  $x = 0.5$ , and finally has a maximum ( $= 6\%$ ) for  $x = 0.75$ . As illustrated in Figure 1c,  $d_p$  is the density that includes the volume of the closed pores in the particles. Hence, the large difference in the values for  $d_{\text{XRD}}$  and  $d_p$  for the samples with  $x > 0.5$  indicates the presence of closed pores in the particles. A detailed analysis for such difference is presented in the Section 3.4.

Figure 7 shows the values of (a)  $Q_{\text{vol}}$ , (b)  $W_{\text{vol}}$ , and (c)  $W_{\text{vol}}(d_p)/W_{\text{vol}}(d_{\text{XRD}})$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples, where the



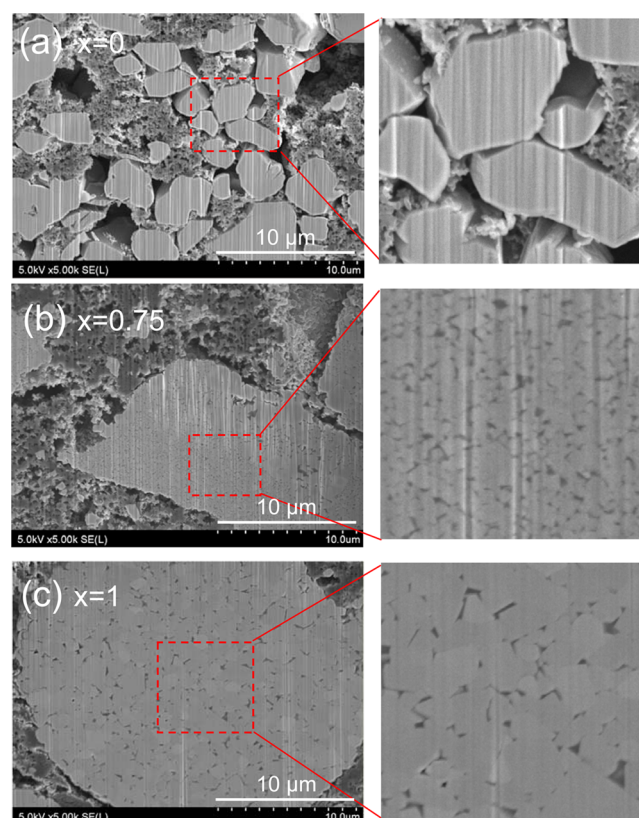
**Figure 7.** (a) Volumetric discharge capacity ( $Q_{\text{vol}}$ ), (b) volumetric energy density ( $W_{\text{vol}}$ ), and (c) the ratio of  $W_{\text{vol}}(d_p)/W_{\text{vol}}(d_{\text{XRD}})$  as a function of  $x$  in  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ . Note that the  $W_{\text{vol}}(d_p)/W_{\text{vol}}(d_{\text{XRD}})$  ratio is exactly same to the  $Q_{\text{vol}}(d_p)/Q_{\text{vol}}(d_{\text{XRD}})$  ratio, because the  $W_{\text{vol}}(d_p)$  and  $W_{\text{vol}}(d_{\text{XRD}})$  values are calculated by  $W_{\text{vol}}(d_p) = Q_{\text{vol}}(d_p) V_{\text{ave}}$  and  $W_{\text{vol}}(d_{\text{XRD}}) = Q_{\text{vol}}(d_{\text{XRD}}) V_{\text{ave}}$ , respectively.

$Q_{\text{vol}}$  and  $W_{\text{vol}}$  were calculated using both  $d_{\text{XRD}}$  and  $d_p$  values. The  $x$  dependence of  $Q_{\text{vol}}$  calculated using  $d_{\text{XRD}}$  is similar to that of  $Q_{\text{gra}}$  (see Figure 5b); as  $x$  increases from 0,  $Q_{\text{vol}}$  slightly decreases in region (i), then increases with a change in the slope ( $dQ_{\text{vol}}/dx$ ) in region (ii), and finally has a maximum at  $x = 0.9$  in region (iii). The maximum and minimum  $Q_{\text{vol}}$  values are  $894 \text{ mAh cm}^{-3}$  for  $x = 0.9$  and  $652 \text{ mAh cm}^{-3}$  for  $x = 0.25$ , respectively. At these compositions, the  $Q_{\text{gra}}$  values also have the maximum and minimum values, i.e.,  $186 \text{ mAh g}^{-1}$  for  $x = 0.9$  and  $133 \text{ mAh g}^{-1}$  for  $x = 0.25$  (Figure 5a and Table 1). However, the  $Q_{\text{vol}}(\text{min})/Q_{\text{vol}}(\text{max})$  ratio is smaller than the  $Q_{\text{gra}}(\text{min})/Q_{\text{gra}}(\text{max})$  ratio, because  $d_{\text{XRD}}$  monotonically

decreases with  $x$ . For the  $Q_{\text{vol}}$  values calculated using  $d_p$ , the differences between the samples are smaller because of the large decrease in the values for  $d_p$  compared to those for  $d_{\text{XRD}}$ , particularly for the samples with  $x > 0.5$ . As a result, the maximum and minimum  $Q_{\text{vol}}$  values obtained using  $d_p$  are  $858 \text{ mAh cm}^{-3}$  for  $x = 0.9$  and  $652 \text{ mAh cm}^{-3}$  for  $x = 0.25$ , respectively.

As can be seen in Figure 7b, the  $x$  dependence of  $W_{\text{vol}}$  is essentially similar to that of  $Q_{\text{vol}}$ . The value for  $W_{\text{vol}}$  obtained using  $d_{\text{XRD}}$  ( $d_p$ ) decreases from  $2764$  ( $2720$ )  $\text{mWh cm}^{-3}$  at  $x = 0$  to  $2540$  ( $2479$ )  $\text{mWh cm}^{-3}$  at  $x = 0.25$ , then increases to  $2945$  ( $2797$ )  $\text{mWh cm}^{-3}$  at  $x = 0.667$ , and finally reaches a maximum of  $3407$  ( $3268$ )  $\text{mWh cm}^{-3}$  at  $x = 0.9$ . It should be noted that the difference in the values of  $W_{\text{vol}}$  calculated using  $d_{\text{XRD}}$  and  $d_p$  is  $\sim 1.6\%$  at  $x = 0$  and  $\sim 4.3\%$  at  $x = 0.75$ . This variation suggests that  $W_{\text{vol}}^{\text{act}}(Q_{\text{vol}}^{\text{act}})$  for the  $x = 0.75$  sample is less than that for the  $x = 0$  sample, because of the larger difference in the values for  $d_{\text{XRD}}$  and  $d_p$ . For the same reason, the  $W_{\text{vol}}(d_p)/W_{\text{vol}}(d_{\text{XRD}})$  ratio has the minimum at the  $x = 0.75$  composition (see Figure 7c). Note that the  $W_{\text{vol}}(d_p)/W_{\text{vol}}(d_{\text{XRD}})$  ratio is exactly same to the  $Q_{\text{vol}}(d_p)/Q_{\text{vol}}(d_{\text{XRD}})$  ratio, because the  $W_{\text{vol}}(d_p)$  and  $W_{\text{vol}}(d_{\text{XRD}})$  values are calculated by  $W_{\text{vol}}(d_p) = Q_{\text{vol}}(d_p) \times V_{\text{ave}}$  and  $W_{\text{vol}}(d_{\text{XRD}}) = Q_{\text{vol}}(d_{\text{XRD}}) \times V_{\text{ave}}$ , respectively.

**3.4. Cross-Sectional Observations.** To clarify the origin of the significantly smaller values for  $d_p$  as compared to  $d_{\text{XRD}}$ , particularly when  $x > 0.5$ , cross-sectional views of the samples were obtained, and the results are shown in Figure 8: (a)  $x = 0$ , (b)  $x = 0.75$ , and (c)  $x = 1$ . The primary particles in the  $x = 0$



**Figure 8.** Cross-sectional views for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with (a)  $x = 0$ , (b)  $x = 0.75$ , and (c)  $x = 1$ . Enlarged cross-sectional views are also shown to clarify the existence/absence of closed pores and grain boundaries in the particles.

sample are separated from one another, and their average size is approximately 3  $\mu\text{m}$ . In addition, neither grain boundaries nor closed pores are observed in these primary particles, as can be seen in the enlarged cross-sectional view of Figure 8a. On the contrary, for the  $x = 0.75$  sample, the primary particles located within  $\sim 0.3 \mu\text{m}$  are agglomerated together, forming large secondary particles of  $\sim 20 \mu\text{m}$ . Furthermore, as can be clearly seen in the enlarged cross-sectional view of Figure 8b, many closed pores exist in the secondary particles. Thus, the significantly smaller value for  $d_p$  compared to that of  $d_{\text{XRD}}$  can be attributed to the presence of closed pores in the secondary particles, because the value of  $d_p$  does not include these volumes (Figure 1c). For the  $x = 1$  sample, the situation is similar to that for the  $x = 0.75$  sample; however, the size of the closed pores appears to be larger than that for the  $x = 0.75$  sample (see the enlarged cross-sectional view of Figure 8c). This suggests that the average size of the primary particles for the  $x = 1$  sample ( $\sim 1 \mu\text{m}$ ) is larger than that for the  $x = 0.75$  sample.

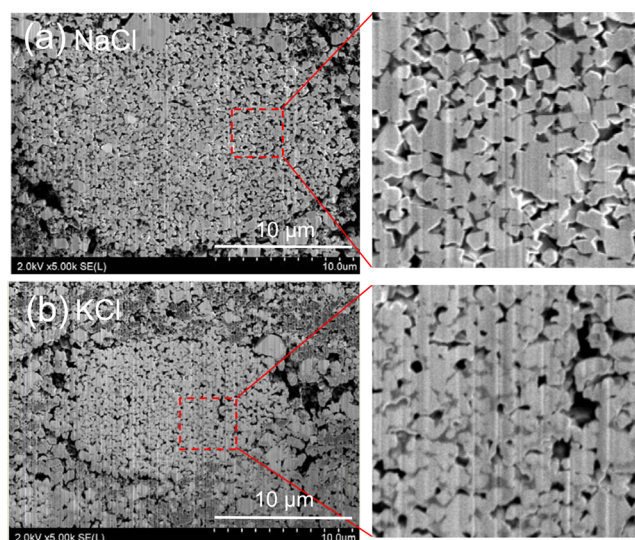
The presence of closed pores unambiguously reduces the value of  $W_{\text{vol}}^{\text{act}}$  ( $Q_{\text{vol}}^{\text{act}}$ ) for  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ . Moreover, it leads to possible capacity fading, especially for high-rate charge and discharge tests, because electrical conduction is only achieved from the surfaces of secondary particles. That is, if each primary particle is isolated from other primary particles due to a change in  $a_h$ ,  $c_h$ , or  $V_h$ , the electrical isolation of the active primary particles occurs. The advantages of cross-sectional observations using FIB rather than conventional SEM analyses should also be emphasized. Although our recent SEM observations of stoichiometric  $\text{LiCoO}_2$  indicated agglomerated particle morphology,<sup>7</sup> the cross-sectional views shown in Figure 8a clarified that the primary particles for the  $x = 0$  sample are distinct from one another. It should be also noted that that conducting carbon and binder are added when forming mixed electrodes. Thus, the observations of not only active materials but also mixed electrodes are essential for elucidating the distribution of the active materials.

The agglomerated particle morphology described above is characteristic of Ni-rich  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  compounds, as reported for  $\text{Li}(\text{Co}_{0.1}\text{Ni}_{0.8}\text{Mn}_{0.1})\text{O}_2$ <sup>21</sup> and  $\text{LiCo}_{0.15}\text{Ni}_{0.82}\text{Al}_{0.03}\text{O}_2$ .<sup>25</sup> With respect to the sample with  $x = 0$  ( $\text{LiCoO}_2$ ) and Ni-rich  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  compounds, there are several possibilities for the origin of the morphological differences. One of the most probable factors is the synthesis temperature.  $\text{LiCoO}_2$  is typically prepared at 900  $^\circ\text{C}$  in order to obtain a highly crystallized material. However, for Ni-rich  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  compounds, high temperature synthesis above 750  $^\circ\text{C}$  inevitably induces the formation of a rock-salt phase  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  ( $Fm\bar{3}m$ ) with  $z > \sim 0.1$ , which is electrochemically inactive.<sup>9–11</sup> Thus, the synthesis temperature for Ni-rich  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  materials is typically limited to  $\sim 750 \text{ }^\circ\text{C}$ ,<sup>9,13,14</sup> and consequently, is not sufficient for obtaining well-developed primary particles. In this study, a synthesis temperature of 750  $^\circ\text{C}$  was employed for the samples with  $x \geq 0.75$ . The nearly linear  $x$  dependence in  $d_p$  can also be explained by the synthesis temperature. A second factor is the surface energy.  $\text{LiCoO}_2$ , particularly when prepared in a flux medium,<sup>26</sup> exhibits a hexagonal shape because of the anisotropic crystal growth along the  $[00\bar{1}]$  direction, whereas  $\text{LiNiO}_2$  and its derivatives exhibit nonuniform particle morphologies.<sup>21,25</sup> As pointed out by Kim and co-workers,<sup>21,25</sup> this difference implies that the surface energy of  $\text{LiNiO}_2$  is less than that of  $\text{LiCoO}_2$ . Further theoretical studies on surface energy on  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$

with  $x$  would provide useful information for determining the particle morphology of  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$ .

**3.5. Attempt to Increasing  $d_p$ .** The cross-sectional observations of the electrodes revealed that significantly smaller value for  $d_p$  as compared with  $d_{\text{XRD}}$ , particularly for the samples with  $x > 0.5$ , was due to the presence of closed pores in agglomerated secondary particles. Therefore, an attempt was made to increase  $d_p$  for the  $x = 0.75$  sample, because the value of  $d_{\text{XRD}}/d_p$  reached the maximum for this composition (Figure 6b). One approach for obtaining high  $d_p$  values is to synthesize well-developed primary particles, such as those observed for the  $x = 0$  sample (Figure 8a). For this purpose, the flux method using NaCl or KCl as the medium was employed, as reported for  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$ ,<sup>21</sup>  $\text{LiCoO}_2$ ,<sup>26</sup> and  $\text{Li}[\text{Li}_{1/3}\text{Ti}_{5/3}]\text{O}_4$ .<sup>27</sup> Such compounds were typically heated at 900–1000  $^\circ\text{C}$ ,<sup>21,26,27</sup> because the melting points of NaCl and KCl are 800 and 770  $^\circ\text{C}$ , respectively. However, in this study, the heating temperature was restricted to 750  $^\circ\text{C}$  in order to minimize the value of  $z$  in  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$ .

As expected, the  $x = 0.75$  samples prepared using NaCl or KCl as a flux medium have greater  $d_p$  values compared to that for the conventional  $x = 0.75$  sample;  $d_p = 4.85(4)$  and  $4.78(4) \text{ g cm}^{-3}$  for NaCl and KCl, respectively, while  $d_p = 4.60(2) \text{ g cm}^{-3}$  for the conventional  $x = 0.75$  sample. These values are very close to the  $d_{\text{XRD}}$  value ( $= 4.871(1) \text{ g cm}^{-3}$ ) for the conventional  $x = 0.75$  sample (Figure 6a and Table 1). Figure 9

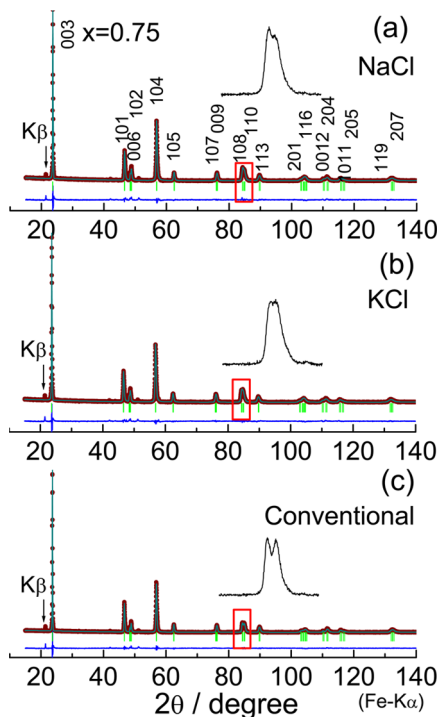


**Figure 9.** Cross-sectional views for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $x = 0.75$  prepared in (a) NaCl and (b) KCl media. The samples were synthesized by heating the reaction mixtures at 750  $^\circ\text{C}$  in air for 18 h. Enlarged cross-sectional views are also shown to clarify the existence/absence of closed pores and grain boundaries in the particles.

shows the cross-sectional views for the  $x = 0.75$  samples prepared using (a) NaCl and (b) KCl as a flux medium. The primary particles prepared using flux media, with an average size of  $\sim 0.5 \mu\text{m}$ , are larger than those for the conventional  $x = 0.75$  sample ( $< \sim 0.3 \mu\text{m}$ ). However, the primary particles are not as developed as those in the  $x = 0$  sample (Figure 8a), and are still agglomerated, forming large secondary particles ( $\sim 25 \mu\text{m}$  for NaCl and  $\sim 20 \mu\text{m}$  for KCl). It can be also clearly seen in Figures 9a and 9b that each type of primary particles is distributed in a more separate form, compared to those of the conventional  $x = 0.75$  and 1 samples (Figure 8b, c,

respectively). Thus, opened, rather than closed, pores in the secondary particles are considered to contribute to higher  $d_p$  values.

Next, we performed the XRD measurements on the  $x = 0.75$  samples prepared using the NaCl and KCl flux media by iron  $K\alpha$  radiation. Figure 10 shows the results of Rietveld analyses



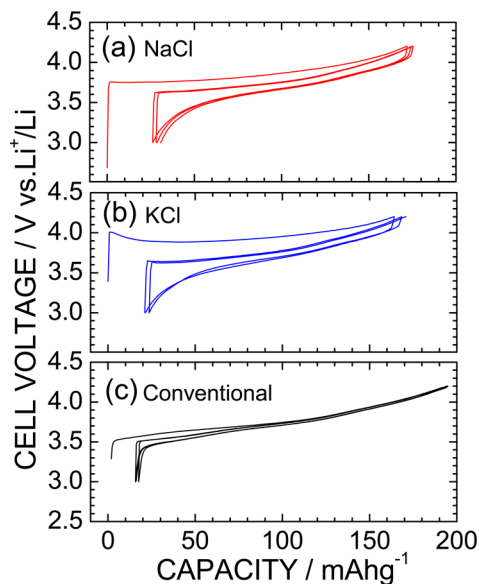
**Figure 10.** Results of the Rietveld analyses for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $x = 0.75$  prepared in (a) NaCl and (b) KCl media. (c) Result for the conventional  $x = 0.75$  sample shown for comparison. Enlarged XRD patterns at approximately  $2\theta = 85^\circ$  are shown in the inset. The diffraction line indicated by arrow is due to the iron  $K\beta$  radiation.

for the samples prepared using (a) NaCl and (b) KCl as flux media. The result of Rietveld analysis for the conventional  $x = 0.75$  sample is also shown in c for comparison. The crystal structure of both a and b is assigned as a single-phase with layered structure with the  $R\bar{3}m$  space group. Impurity phases, including  $\text{Co}_3\text{O}_4$  and/or  $\text{NiO}$ , are barely detected in both XRD patterns, where  $d_{\text{XRD}} = 6.11 \text{ g cm}^{-3}$  for  $\text{Co}_3\text{O}_4$  and  $d_{\text{XRD}} = 6.67 \text{ g cm}^{-3}$  for  $\text{NiO}$ . The lattice parameters are calculated as  $a_h = 2.871(1) \text{ \AA}$  and  $c_h = 14.15(1) \text{ \AA}$  for (a), and  $a_h = 2.871(1) \text{ \AA}$  and  $c_h = 14.16(1) \text{ \AA}$  for (b), which are similar to those for the conventional  $x = 0.75$  sample (Table 1). However, the integrated peak ratio for the 003 and 104 diffraction lines  $I(003)/I(104)$  is clearly different between the samples prepared using the flux media and the conventional sample. The  $I(003)/I(104)$  values are 1.00, 1.02, and 1.72 for the NaCl, KCl, and conventional  $x = 0.75$  samples. This result suggests that, as reported previously,<sup>9–14</sup> a large amount of  $z$  exists in  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  samples prepared by the flux method. Indeed, the Rietveld analyses revealed that  $z = 0.06(1)$  for NaCl and  $z = 0.05(1)$  for KCl, but  $z \approx 0$  for the conventional  $x = 0.75$  sample. The enlarged XRD patterns at approximately  $2\theta = 85^\circ$  also provide another insight into the crystal structure. For the samples prepared using NaCl and KCl media, the 108 diffraction line is located near the 110

diffraction line, while, for the conventional  $x = 0.75$  sample, the 108 and 110 diffraction lines are clearly separated from one another. According to the structural relationship between cubic and hexagonal systems, the 440 diffraction line of the cubic system is converted to the 108 and 110 diffraction lines of the hexagonal system, when  $c_h/a_h \neq 2\sqrt{6}$  ( $\sim 4.90$ ).<sup>12</sup> Therefore, it can be concluded that the synthesis by the flux method using NaCl or KCl medium induces a structural change from the hexagonal phase to the cubic phase, simultaneously, although such method provides high  $d_p$  values at the  $x = 0.75$  composition.

It should be noted that both phenomena (high  $d_p$  value and structural change) are also observed for the  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$  compound, which was prepared using the NaCl flux medium at  $900^\circ\text{C}$ .<sup>21</sup> That is, the  $d_p$  value of such compound is  $4.83 \text{ g cm}^{-3}$ ,<sup>21</sup> which is slightly higher than the expected  $d_{\text{XRD}}$  value ( $\sim 4.81 \text{ g cm}^{-3}$  or less, because the substitution of Mn ions for Co ions decreases  $M_w$ ). Moreover, the XRD measurements indicated the cubic structural character; the peak intensity of the 003 diffraction line is almost similar to that of the 104 diffraction line, and the 108 and 110 diffraction lines are not clearly separated.<sup>21</sup>

Considering from the previous results on the  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  compounds,<sup>9–14</sup> the large amount of  $z$  and cubic structural character significantly influence the electrochemical reactivity. We, thus, examined the electrochemical reactivity of the  $x = 0.75$  samples prepared using the NaCl and KCl flux media. Figure 11 shows the charge



**Figure 11.** Charge and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with  $x = 0.75$  prepared in (a) NaCl and (b) KCl media. (c) Charge and discharge curves for the conventional  $x = 0.75$  sample shown for comparison.

and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with  $x = 0.75$  prepared in the (a) NaCl and (b) KCl media. The charge and discharge curves for both samples significantly differ from those for the conventional  $x = 0.75$  sample (see Figure 10c). For instance, the  $V_{\text{cell}}$  values at the beginning of the charge reaction rapidly increase to  $\sim 3.8 \text{ V}$  for NaCl and  $\sim 4.0 \text{ V}$  for KCl, while that is  $\sim 3.5 \text{ V}$  for the conventional  $x = 0.75$  sample. It is empirically known that the charge and discharge curves at the initial cycle are different from those at the subsequent



cycles, due to the decomposition reaction between active material (or lithium metal) and electrolyte.<sup>28</sup> Here the samples prepared using the NaCl and KCl flux media were washed with distilled water after the synthesis. Thus, the rapid increase in  $V_{\text{cell}}$  at the beginning of the charge reaction is probably attributed to the decomposition reaction between the adsorbed water on the sample and electrolyte, as in the case for the chemically delithiated  $\text{Li}_x\text{NiO}_2$  compound.<sup>29</sup> Besides of the difference at the initial cycle, the hysteresis between the charge and discharge curves, i.e., the polarization for the NaCl and KCl samples is larger than that for the conventional  $x = 0.75$  sample. The polarization for both samples is  $\sim 0.4$  V at the middle point of the charge and discharge curves, whereas that for the conventional  $x = 0.75$  sample is less than 0.05 V. This is caused by the difference in crystal structure between the samples, as clarified by the Rietveld analyses; the large amount of  $z$  hinders the diffusion of the  $\text{Li}^+$  ions, because the rock-salt phase  $(\text{Li}_{1-z}\text{Ni}_z)_{3b}[\text{Co}_{1-x}\text{Ni}_{x-z}]_{3a}\text{O}_2$  ( $Fm\bar{3}m$ ) with  $z > \sim 0.1$  is electrochemically inactive. As a result, the  $Q_{\text{gra}}$  values are 146 mAh  $\text{g}^{-1}$  for NaCl and 145 mAh  $\text{g}^{-1}$  for KCl, which are  $\sim 33$  mAh  $\text{g}^{-1}$  less than that for the conventional  $x = 0.75$  sample.

As described above, we have a dilemma concerning high  $d_p$  and high  $W_{\text{vol}}$  values on the  $x = 0.75$  sample. The actual  $W_{\text{vol}}$  value for the conventional  $x = 0.75$  sample was decreased, due to the smaller value for  $d_p$  ( $\sim 4.6$  g  $\text{cm}^{-3}$ ) as compared with  $d_{\text{XRD}}$  [ $= 4.871(1)$  g  $\text{cm}^{-3}$ ]. The high  $d_p$  values ( $\sim 4.8$  g  $\text{cm}^{-3}$ ) were obtained by the flux method, but the flux method also induced the structural changes on  $z$  and  $I(003)/I(104)$ . Therefore, the  $Q_{\text{gra}}$  values, i.e., the  $W_{\text{vol}}$  values for the samples prepared with the flux method were less than that for the conventional  $x = 0.75$  sample. The situation for the  $x > 0.75$  samples is thought to be similar to that for the  $x = 0.75$  sample, because the amount of  $z$  increases with  $x$  (section 3.1). Although the optimum synthetic condition for obtaining both high  $d_p$  and high  $W_{\text{vol}}$  values is currently unknown, the results on the  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$  compounds<sup>21</sup> would provide crucial information to overcome such dilemma. As described above, the  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$  compound prepared with the NaCl flux medium indicated high  $d_p$  value close to  $d_{\text{XRD}}$  and the cubic structural character in terms of the  $I(003)/I(104)$  value. Nevertheless, the charge and discharge curves for such compound are almost similar to those for the conventional  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$  compound, which has almost perfect layered structure.<sup>21</sup> Although the origin of the difference behavior of  $\text{LiCo}_{0.25}\text{Ni}_{0.75}\text{O}_2$  and  $\text{Li}(\text{Ni}_{0.8}\text{Co}_{0.1}\text{Mn}_{0.1})\text{O}_2$  is currently unclear, it is expected that the existence of the Mn ions in the Li layer should affect their electrochemical properties. This is because the  $Q_{\text{gra}}$  value for the  $\text{LiNi}_{1/2}\text{Mn}_{1/2}\text{O}_2$  compound is  $\sim 200$  mAh/g, although its  $I(003)/I(104)$  value indicates the presence of the Mn ions in the Li layer.<sup>15</sup> The present study focused on the  $x$  dependences of  $d_p$ ,  $d_{\text{XRD}}$ ,  $W_{\text{vol}}$  for the solid solution samples  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  with  $0 \leq x \leq 1$ . Further density measurements of  $\text{LiCo}_{1-x-y}\text{Ni}_x\text{Mn}_y\text{O}_2$  associated their  $Q_{\text{gra}}$ ,  $Q_{\text{vol}}$ , and  $W_{\text{vol}}$  would offer essential information for increasing actual  $W_{\text{vol}}$  of LIBs.

## 4. CONCLUSION

The volumetric energy density  $W_{\text{vol}}$  is one of the most important parameters for high-energy density lithium-ion batteries (LIBs). A systematic study of  $Q_{\text{gra}}$ ,  $Q_{\text{vol}}$ ,  $W_{\text{vol}}$ ,  $d_p$ , and  $d_{\text{XRD}}$  for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples with  $0 \leq x \leq 1$  was performed in order to clarify the factors affecting the  $W_{\text{vol}}$ . The

$x$  dependence of  $Q_{\text{gra}}$ ,  $Q_{\text{vol}}$ , and  $W_{\text{vol}}$  was divided into three distinct regions. For instance, the value of  $W_{\text{vol}}$  calculated using  $d_p$  slightly decreased from 2720 mWh  $\text{cm}^{-3}$  for  $x = 0$  to 2479 mWh  $\text{cm}^{-3}$  for  $x = 0.25$ , then rapidly increased to 2797 mWh  $\text{cm}^{-3}$  for  $x = 0.667$  with a change in the slope ( $dW_{\text{vol}}/dx$ ), and finally reached a maximum ( $= 3268$  mWh  $\text{cm}^{-3}$ ) for  $x = 0.9$ . Because the structural parameters, including  $a_b$ ,  $c_b$ , and  $V_b$ , decreased linearly with  $x$ ,  $d_{\text{XRD}}$  monotonically decreased from 5.062(1) g  $\text{cm}^{-3}$  for  $x = 0$  to 4.779(1) g  $\text{cm}^{-3}$  for  $x = 1$ . On the contrary, the  $x$  dependence of  $d_p$  was not monotonic. That is,  $d_p$  decreased almost linearly in the  $x$  range between 0 and 0.5, then  $d_p$  rapidly dropped to 4.63(2) g  $\text{cm}^{-3}$  at  $x > 0.5$ , and finally leveled off to a constant value ( $\sim 4.60$  g  $\text{cm}^{-3}$ ) with further increasing  $x$ . Cross-sectional observations revealed that significantly lower values for  $d_p$  as compared to  $d_{\text{XRD}}$ , particularly for  $x > 0.5$ , is attributed to the presence of closed pores in agglomerated secondary particles. Therefore, the actual  $W_{\text{vol}}$  decreases from the ideal  $W_{\text{vol}}$ , particularly for  $x > 0.5$ , because of the presence of closed pores in secondary particles. In other words, the synthesis of well-developed primary particles without closed pores is essential for increasing the actual  $W_{\text{vol}}$ . The synthesis of a sample with well-developed primary particles was thus attempted using NaCl or KCl as a flux medium. Although high  $d_p$  values were obtained for the  $x = 0.75$  samples prepared in this manner, a structural change from the hexagonal phase to cubic phase was also induced. Hence, the values of  $Q_{\text{gra}}$  for such samples were less than that for the conventional sample. Although it seems that current LIBs rely on electrode packing technologies, these results suggest that a deeper understanding of the differences between  $d_{\text{ele}}$  and  $d_p$  (or  $d_{\text{XRD}}$ ) would offer crucial information for the development of high-energy density LIBs.

## ■ ASSOCIATED CONTENT

### Supporting Information

Schematic illustration of the gas pycnometer system,  $d_p$  for the  $x = 0$  sample as a function of each measurement, XRD patterns for the  $x = 0.05, 0.1, 0.333, 0.667, 0.9$ , and  $0.95$  samples, the Rietveld analyses for the  $x = 0.9, 0.95$ , and  $1$  samples, charge and discharge curves for the  $\text{Li}/\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  cells with  $x = 0.05, 0.1, 0.333, 0.667, 0.9$ , and  $0.95$ , and the average operating voltage for the  $\text{LiCo}_{1-x}\text{Ni}_x\text{O}_2$  samples. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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